## ON BLOW-UP AT SPACE INFINITY FOR SEMILINEAR HEAT EQUATIONS

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We are interested in solutions of semilinear heat equations which blow up at space infinity.

In [7], we considered a nonnegative blowing up solution of

$$u_t = \Delta u + u^p, \qquad x \in \mathbb{R}^n, \qquad t > 0$$

with initial data  $u_0$  satisfying

$$0 \le u_0(x) \le M$$
,  $u_0 \not\equiv M$  and  $\lim_{|x| \to \infty} u_0(x) = M$ ,

where p > 1 and M > 0 is a constant. We proved in [7] that the solution u blows up exactly at the blow-up time for the spatially constant solution with initial data M. We moreover proved that u blows up only at the space infinity. In this paper we would like to generalize this result in the following directions.

- (i) (Initial data) We consider more general initial data  $u_0$  which may not converge to M for all directions of x, for example  $u_0 \to M$  as  $|x| \to \infty$  only for x in some sector. It is convenient to introduce a notion of blow up direction at the space infinity. We are able to give necessary and sufficient condition so that a particular direction is a blow-up direction.
- (ii) (Nonlinear term) We extend the class of nonlinearities. It includes  $e^u$  and  $u^p + u^q$  for p, q > 1.
- In [8] we consider solutions of the initial value problem for the equation

(1) 
$$\begin{cases} u_t = \Delta u + f(u), \quad x \in \mathbb{R}^n, \quad t > 0, \\ u(x,0) = u_0(x), \quad x \in \mathbb{R}^n. \end{cases}$$

The nonlinear term f is assumed to be nonnegative and locally Lipschitz in  $\mathbb{R}$  with the property that

(2) 
$$\lim_{b \ge b_0, \, \delta \in (\delta_0, 1)} \frac{\delta^p f(b)}{f(\delta b)} > 0 \quad \text{for} \quad b_0 > 0, \quad \delta_0 \in (0, 1), \quad p > 1.$$

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We take two constants M and N satisfying M + N > 0 and

$$(3) f(M) > 0.$$

The initial data  $u_0$  is assumed to be a measureable function in  $\mathbb{R}^n$  satisfying

(4) 
$$-N \le u_0 \le M$$
 a.e. and  $u_0 \not\equiv M$  a.e.

We are interested in initial data such that  $u_0 \to M$  as  $|x| \to \infty$  for x in some sector of  $\mathbb{R}^n$ . We assume that

(5) 
$$\operatorname{essinf}_{x \in \tilde{B}_m}(u_0(x) - M_m(x - x_m)) \ge 0 \quad \text{for } m = 1, 2, \dots,$$

where

(6) 
$$\tilde{B}_m = B_{r_m}(x_m)$$

with a sequence  $\{r_m\}$  and a sequence of vectors  $\{x_m\}_{m=1}^{\infty}$  and a sequence of functions  $\{M_m(x)\}$  satisfying

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$$\lim_{m \to \infty} r_m = \infty, \qquad M_m(x) \le M_{m+1}(x) \qquad \text{for } m \ge$$
$$\lim_{m \to \infty} \inf_{s \in [1, r_m]} \frac{1}{|B_s|} \int_{B_s(0)} M_m(x) \, \mathrm{d}x = M.$$

Here  $B_r(x)$  denotes the closed ball of radius r centered at x. (In fact, it follows from (4) that  $|x_m| \to \infty$  as  $m \to \infty$ .)

Problem (1) has a unique bounded solution at least locally in time. However, the solution may blow up in finite time. For a given initial value  $u_0$  and nonlinear term f let  $T^* = T^*(u_0, f)$  be the maximal existence time of the solution. If  $T^* = \infty$ , the solution exists globally in time. If  $T^* < \infty$ , we say that the solution blows up in finite time. It is well known that

(7) 
$$\limsup_{t \to T^*} \|u(\cdot, t)\|_{\infty} = \infty,$$

where  $\|\cdot\|_{\infty}$  denotes the  $L^{\infty}$ -norm in space variables.

In this paper, we are interested in the behavior of a blowing up solution near space infinity as well as the location of blow-up points defined below. A point  $x_{BU} \in \mathbb{R}^n$  is called a *blow-up point* (with value  $\pm \infty$ ) if there exists a sequence  $\{(x_m, t_m)\}_{m=1}^{\infty}$  such that

 $t_m \uparrow T^*, \quad x_m \to x_{BU} \quad \text{and} \quad u(x_m, t_m) \to \pm \infty \quad \text{as} \quad m \to \infty.$ 

If there exists a sequence  $\{(x_m, t_m)\}_{m=1}^{\infty}$  such that

 $t_m \uparrow T^*, \qquad |x_m| \to \infty \qquad \text{and} \qquad u(x_m, t_m) \to \pm \infty \quad \text{as} \quad m \to \infty,$ 

then we we say that the solution blows up to  $\pm \infty$  at space infinity.

A direction  $\psi \in S^{n-1}$  is called a blow-up direction for the value  $\pm \infty$  if there exists a sequence  $\{(x_m, t_m)\}_{m=1}^{\infty}$  with  $x_m \in \mathbb{R}^n$  and  $t_m \in (0, T^*)$  such that  $u(x_m, t_m) \to \pm \infty$  (as  $m \to \infty$ ) and

(8) 
$$\frac{x_m}{|x_m|} \to \psi$$
 as  $m \to \infty$ .

We consider the solution v(t) of an ordinary differential equation

(9) 
$$\begin{cases} v_t = f(v), \quad t > 0, \\ v(0) = M. \end{cases}$$

Let  $T_v = T^*(M, f)$  be the maximal existence time of the solution of (9), i. e.,

$$T_v = \int_M^\infty \frac{\mathrm{d}s}{f(s)}.$$

We are now in position to state our main results.

**Theorem 1.** Assume that f is locally Lipschitz in  $\mathbb{R}$  and satisfies (2) and (3). Let  $u_0$  be a continuous function satisfying (4) and (5), and  $T_v \leq T^*(-N, f)$ . Then there exists a subsequence of  $\{x_m\}_{m=1}^{\infty}$  (still denoted by  $\{x_m\}$ , independent of t) such that

$$\lim_{m \to \infty} u(x_m, t) = v(t).$$

The convergence is uniform in every compact subset of  $\{t : 0 \le t < T_v\}$ . Moreover, the solution blows up at  $T_v$ .

**Remark.** Our assumption  $T_v \leq T^*(-N, f)$  says that the solution does not blow up to minus infinity before it blows up to plus infinity. From the condition (4), it follows that  $\lim_{m\to\infty} |x_m| = \infty$ .

This result in particular implies that

(10) 
$$\sup_{0 < t < T^*} v^{-1}(t) \| u(\cdot, t) \|_{\infty} < \infty.$$

When we set  $f(u) = |u|^{p-1}u$ , such a blow-up rate estimate is known for subcritical p; see e.g. [3], [5], [6] for general bounded initial data without assuming (4) and (5). However, for supercritical p such a blow-up rate estimate (10) may not hold in general; see e.g. [1], [9]. If one considers only radial solutions of (1) for supercritical p less than  $1 + 4/(n - 4 - 2(n - 1)^{1/2})$  or  $n \leq 10$ , then the estimate (10) holds [11]. We would like to emphasize that Theorem 1 does not require any restriction on p.

Our second main result is on the location of blow-up points.

**Theorem 2.** Assume the same hypotheses as in Theorem 1. Then the solution of (1) has no blow-up points with  $+\infty$  in  $\mathbb{R}^n$ . (It blows up only at space infinity.)

There is a huge literature on location of blow-up points since the work of Weissler [13] and Friedman-McLeod [2]. (We do not intend to list references exhaustively in this paper.) However, most results consider either bounded domains or solutions decaying at space infinity; such a solution does not blow up at space infinity [4].

As far as the authors know, before the result of [7] the only paper discussing blow-up at space infinity is the work of Lacey [10]. He considered the Dirichlet problem in a half line. He studied various nonlinear terms and proved that a solution blows up only at space infinity. In particular, his result implies that the solution of

$$\begin{cases} u_t = u_{xx} + f(u), & x > 0, \ t > 0, \\ u(0,t) = 1, & t > 0, \\ u(x,0) = u_0(x) \ge 1, \ x > 0 \end{cases}$$

blows up only at space infinity, where  $u_0$  satisfies  $0 \le u_0 \le M$  with M > 1, and  $f(s) = s^p$  and  $e^s$ .

His method is based on construction of suitable subsolutions and supersolutions. However, the construction heavily depends on the Dirichlet condition at x = 0 and does not apply to the Cauchy problem even for the case n = 1.

As previously described, the authors [7] proved the statement of Theorems 1 and 2 assuming that  $u_0(x) \leq M$  for sufficiently large M for positive solutions of  $u_t = \Delta u + u^p$ . Later, Shimojyo [12] had the same results as in [7] by relaxing the assumptions of initial data  $u_0 \geq 0$  which is similar to that in the present paper. His approach is a construction of a suitable supersolution which implies that  $a \in \mathbb{R}^n$  is not a blow-up point. Although he restricted himself to  $f(s) = s^p$ , his idea works for our f under slightly stronger assumption on  $u_0$ . Here we give a different approach.

From Shimojyo's results [12], there arises a problem of "blow-up direction" defined in (8). We next study this "blow-up direction" for the value  $+\infty$ . Our third result is on this blow-up direction. It is convenient to introduce the function  $A_m$  defined by

(11) 
$$A_m(s) = \frac{1}{|B_s(y_m)|} \int_{B_s(y_m)} u_0(z) \, \mathrm{d}z$$

for a given sequence  $\{y_m\}_{m=1}^{\infty}$ . This  $A_m(s)$  represents the mean value of  $u_0$  over the ball  $B_s(y_m)$ .

**Theorem 3.** Assume the same hypotheses as in Theorem 1 and let  $\{s_m\}_{m=1}^{\infty}$  be a sequence diverging to  $\infty$  in  $\mathbb{R}$ . For a given direction  $\psi \in S^{n-1}$ , the following alternatives hold.

(i) If there exists a sequence  $\{y_m\}_{m=1}^{\infty}$  satisfying  $\lim_{m\to\infty} y_m/|y_m| = \psi$  it holds that

$$\limsup_{m \to \infty} \inf_{s \in (1, s_m)} A_m(s) = M,$$

then  $\psi$  is a blow-up direction.

(ii) If for any sequence  $\{y_m\}_{m=1}^{\infty}$  satisfying  $\lim_{m\to\infty} y_m/|y_m| = \psi$  there exists a constant  $c \in (1/(M+N), \infty)$  such that

$$\limsup_{m \to \infty} \inf_{s \in (1,c)} A_m(s) \le M - \frac{1}{c},$$

then  $\psi$  is not a blow-up direction.

This characterizes blow up directions by profiles of initial data. This is a new result even if  $f(u) = |u|^{p-1}u$  or n = 1.

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Here are the main ideas of the proofs. To prove Theorem 1 we construct a suitable subsolution. To prove Theorem 2 we derive a non blow-up criterion. We do not appeal any energy arguments for rescaled function as is done in our previous paper [7]. Our argument consists of two parts. First we observe that

$$u(x,t) \le \delta v(t)$$

near a point  $a \in \mathbb{R}^n$  with some  $\delta \in (0, 1)$  when t is close to the blow-up time. By a bootstrap argument we derive that u is actually bounded near a when t is close to the blow-up time. To prove Theorem 3 we use a comparison argument as in Theorems 1 and 2 and a non blow-up criterion as in the proof of Theorem 2. Moreover, we give conditions on the direction  $\psi \in S^{n-1}$  for being the blow-up direction or not cover all of  $S^{n-1}$  exclusively.

The detailed proofs will be discussed in paper [8].

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